Feature Report

Mixing: Impeller Performance in Stirred Tanks

Characterizing mixer impellers on the basis of power, flow, shear and efficiency

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IN BRIEF TURBULENCE IMPELLER GEOMETRIES HYDRAULIC EFFICIENCY

SHEAR						
PROCESS RESULTS						
TRAILING VORTEX						

APPLICATIONS CONCLUSIONS ixing has been defined as "the application of mechanical motion in order to create fluid dynamic effects that achieve a desired process result" [1]. The process result is the objective of the vessel operator and will be a transformation of the ingredients fed to the vessel into a product. The goal of the equipment supplier will be to understand the role of mixing in promoting the transformation and choosing an impeller that will create the appropriate fluid-dynamic effects to do this.

Processes carried out in stirred tanks can be generally divided into the following two classes:

- Those relying on flow generated by the impeller creating motion throughout the fluid, such as blending of pigments into a resin or emulsion in paint manufacture where homogeneity of the vessel contents is critical to product quality
- Those relying on "shear" to reduce the size of a second dispersed phase, whether gas bubbles, liquid droplets or particles, such as a hydrogenation reactor where smaller bubbles provide more surface area for mass transfer from the gas into the liquid phase

Impellers are often described qualitatively as, among others, high flow, high shear or high efficiency, and the choice of equipment required to achieve the process result most efficiently is made on this vague basis. This article describes how the performance characteristics of impellers commonly used in stirred tanks can be quantified, thereby enabling engineers to make educated decisions about which ones to use in order to achieve their desired process results.

Turbulence

Turbulent flow is characterized by the presence of random fluctuations in velocity, socalled eddies, that are superimposed on the mean, time-averaged flow. There will be a

NOMENCLATURE

- A Constant in Equation (30) A_{DIS} Discharge area for primary flow from impeller
- D Impeller diameter
- d₃₂ Sauter mean droplet size
- FI Flow or pumping number (= $Q/(ND^3)$)
- K Ratio (= $\varepsilon_{MAX}/\overline{\epsilon}$)
- $k_{\ensuremath{\textit{MAX}}}$ Maximum kinetic energy in trailing vortex
- I₀ Diameter of trailing vortex
- N Impeller rotational speed

P Power

- Po Power number [= $P/(\rho N^3 D^5)$]
- Q Flow rate generated by impeller
- R Impeller radius
- *r* Radial position in impeller discharge for estimating velocity gradient
- Re Impeller Reynolds number (= $\rho ND^2/\mu$)
- T Vessel diameter
- \overline{U} Mean velocity in impeller discharge (= Q/A_{DIS})
- V_{TIP} Impeller tip speed

 v_{H} High velocity in impeller discharge for estimating velocity gradient

- $\ensuremath{v_L}$ Low velocity in impeller discharge for estimating velocity gradient
- w Projected blade height
- x Ratio of impeller to trailing vortex diameters (= D/I_0)
- y Distance (in definition of shear rate)
- α Constant (= v_H/V_{TIP})
- β Constant (= v_L/V_{TIP})
- γ Time-averaged velocity gradient
- $ar{arepsilon}$ Power input per mass of fluid in vessel

$$\begin{split} & \varepsilon_{MAX} \ \text{Local energy dissipation rate in trialing vortex} \\ & \eta \ \text{Efficiency defined as mass of fluid pumped per unit} \\ & \text{energy input by impeller} \end{split}$$

- Λ Shear rate constant
- ρ Liquid density
- ϕ Efficiency defined as kinetic energy of fluid divided
- by mechanical energy input by impeller
- Ψ Constant (= r_H/R)
- ω Constant (= r_L/R)

Subscripts

AX Axial HYDFL Hydrofoil IMP Impeller HYDR Hydraulic MECH Mechanical PBT Pitched-blade turbine RD Radial RUSH Rushton



FIGURE 1. Four general classes of impellers are used in stirred tanks operating at low to medium viscosities in the turbulent regime. These impellers primarily generate: axial flow (a, b, c); mixed flow (d), radial flow (e, f, g) or dispersion or de-agglomeration (h)

range of eddy time and length scales associated with a particular flow field. The size of the largest eddies will be on the order of the size of the equipment generating the flow (for example, the blade width of an impeller). The size of the smallest eddies is the Kolmogorov length scale. The eddies also have a lifetime, with the larger eddies existing for a longer period than the small ones. Understanding the role turbulence plays in mixing processes is critical to successful design and scaleup [2].

Impeller geometries

There are four general classes of impellers used in stirred tanks operating in low to medium viscosity fluids in the turbulent regime ($\text{Re} > 10^4$):

1. Axial flow. The primary flow generated by an axial-flow impeller is directed down toward the base of the vessel. Hydrofoils with narrow or wide blades are in this category.

Hydrofoils have profiled blades that may be narrow like an airplane wing (Figure 1a) or wide like a marine propeller (Figure 1b). These impellers were developed to generate the same velocity profile as a propeller, but to be fabricated rather than cast to reduce the impeller's weight and cost. They are also easier to install since they can be supplied as a hub and blades that are assembled inside the vessel [3]. These impellers are generally considered to be "low-shear" [4].

An anti-ragging hydrofoil (Figure 1c) is used in wastewater applications. It has blades that are swept-back preventing build-up of fibrous matter, which is commonly present in municipal wastewater, on the leading edge of the blades.

2. *Mixed flow.* These impellers generate both axial and radial components of velocity and the distribution between the two can be controlled by adjusting the impeller diameter to vessel diameter ratio. Pitched-blade turbines (Figure 1d) are in this category.

Pitched-blade turbines have flat blades that are usually angled at 45 deg, although shallower and steeper angles are sometimes used.

3. Radial flow. These impellers generate a strong radial component of velocity directed at the vessel wall. A pitched blade turbine with 90-deg blade angle generates radial flow and is commonly called a flat-blade turbine (Figure 1e).

Impellers used for processes requiring dispersion of gas bubbles also generate a primarily radial flow, but have blades attached to a disk. The Rushton (Figure 1f) and Smith (Figure 1g) turbines are commonly used for these processes. The disk ensures that bubbles fed into the vessel beneath the impeller must flow through the blades where the local "shear" breaks them up, creating high interfacial area for mass transfer. The Rushton turbine is generally considered to be "high-shear" [4].

4. *High-speed dispersers.* These impellers look like circular-saw blades with alternating teeth angled up and down (Figure 1h). They operate at high rotational and tip speeds and are used almost exclusively for processes that require significant size reduction, such as dispersion and de-agglomeration of dry powder when preparing a slurry from liquid and a dry powder.

Hydraulic efficiency

Impellers in stirred tanks are machines that move fluid; essentially they are pumps. Like pumps, their efficiency can be defined and calculated. The hydraulic efficiency of a pump is the ratio of the kinetic energy of the flowing fluid to the mechanical energy input by the impeller.

The mechanical power input by an impeller in a stirred vessel is calculated from the following equation (Note: all nomenclature are defined in the box on p. 46):

$$P_{MECH} = Po \cdot \rho \cdot N^3 \cdot D^5$$



FIGURE 2. This graph plots the hydraulic efficiency, $\varphi_{\text{HYDR}},$ versus the power number for various impellers

Po is the impeller's power number and it is a drag coefficient that is determined by the geometry of the impeller (blade width, blade angle, number of blades and so on.).

The primary flow generated by an impeller is calculated from Equation (2):

$$Q = Fl \cdot N \cdot D^3 \tag{2}$$

Fl is the impeller's flow, or pumping number.

Both the power and flow numbers are measured experimentally and typical values for commonly used impellers are given in Table 1.

The average velocity in the impeller discharge can be calculated from Equation (3):

$$\overline{U} = \frac{Q}{A_{DIS}} \tag{3}$$

A_{DIS} is the area through which the primary flow is pumped. For axial-flow impellers, this is a disk with diameter equal to the impeller diameter and for radial-flow impellers it is the wall of a cylinder with diameter equal to the impeller diameter and height equal to the blade width. For axial-flow impellers:

$$\overline{U} = \frac{Fl \cdot N \cdot D^3}{\left(\frac{\pi}{4}\right) \cdot D^2} = \frac{4Fl}{\pi} \cdot N \cdot D \tag{4}$$

For radial-flow impellers:

TABLE 1. TYPICAL VALUES OF <i>Po</i> , <i>FI</i> AND <i>x</i> FOR COMMON IMPELLERS								
Impeller	Power number	Flow number	$x = D / I_0$					
Narrow-blade hydrofoil	0.30	0.52	17					
Wide-blade hydrofoil	0.70	0.66						
Pitched-blade turbine	1.50	0.80	16					
Flat-blade turbine	3.00	0.80						
Rushton turbine	5.00	0.65	12					
HSD-Sawtooth	0.10	0.05	12					

$$\overline{U} = \frac{Fl \cdot N \cdot D^3}{\pi \cdot w \cdot D} = \frac{Fl \cdot D}{\pi \cdot w} \cdot N \cdot D$$
(5)

The energy dissipation rate, or power, of the flowing fluid is the product of the flowrate and the head that the pump develops:

$$P_{HYDR} = Q \cdot \Delta H \tag{6}$$

Where:

$$\Delta H = \rho \cdot \frac{\overline{U}^2}{2} \tag{7}$$

Combining Equations (2), (4), (6) and (7), for axial flow impellers:

$$P_{HYDR} = Fl \cdot N \cdot D^3 \frac{\rho}{2} \left(\frac{4Fl}{\pi} \cdot N \cdot D\right)^2 = \frac{8Fl^3}{\pi^2} \cdot \rho \cdot N^3 \cdot D^5 \quad (8)$$

$$\phi_{AX} = \frac{P_{HYDR}}{P_{MECH}} = \frac{8}{\pi^2} \cdot \frac{Fl^3 \cdot \rho \cdot N^3 \cdot D^5}{Po \cdot \rho \cdot N^3 \cdot D^5} = \frac{8}{\pi^2} \cdot \frac{Fl^3}{Po}$$
(9)

Similarly, combining Equations (2), (5), (6) and (7), for radial-flow impellers:

$$P_{HYDR} = Fl \cdot N \cdot D^{3} \cdot \frac{\rho}{2} \cdot \left(\frac{Fl \cdot D}{\pi w} \cdot N \cdot D\right)^{2}$$

$$= \frac{Fl^{3}}{2\pi^{2}} \cdot \left(\frac{D}{w}\right)^{2} \cdot \rho \cdot N^{3} \cdot D^{5}$$
 (10)

$$\phi_{RD} = \frac{P_{HYDR}}{P_{MECH}} = \frac{1}{2\pi^2} \cdot \left(\frac{D}{w}\right)^2 \cdot \frac{Fl^3 \cdot \rho \cdot N^3 \cdot D^5}{Po \cdot \rho \cdot N^3 \cdot D^5} = \frac{1}{2\pi^2} \cdot \left(\frac{D}{w}\right)^2 \cdot \frac{Fl^3}{Po}$$
(11)

The hydraulic efficiency, ϕ , is plotted against the impellers' power numbers in Figure 2. The circular symbols represent data measured by the FMP (Fluid Mixing Processes) consortium [5] using laser-Doppler anemometry and the diamonds represent data measured in the PMSL laboratory using particle-image velocimetry. The data are in agreement showing that measurement technique has no effect on the values of hydraulic efficiency calculated.

The hydrofoils are the most efficient impellers followed by the pitched-blade turbines, then the radial flow flatblade and Rushton turbines. The high-shear disperser impeller is the least efficient, with a hydraulic efficiency of less than 1%. The difference in efficiency within a class of impellers is a result of the impeller to tank diameter ratio. A larger impeller is more efficient and this definition of hydraulic efficiency does not take this into account.

An alternative definition of efficiency has been proposed by Fort and others [6]. This is the mass of fluid pumped per unit of energy input by an impeller:

$$\eta_{HYDR} = \frac{\rho \cdot Q}{P_{MECH}} = \frac{\rho \cdot Fl \cdot N \cdot D^3}{Po \cdot \rho \cdot N^3 \cdot D^5} = \frac{Fl}{Po(N \cdot D)^2} \quad (12)$$

This quantity has units of kilogram of fluid pumped per Joule of energy input by the impeller.

The power input per unit mass of fluid, for a vessel where depth is equal to vessel diameter, can be calculated from:

$$\bar{\varepsilon} = \frac{Po \cdot \rho \cdot N^3 \cdot D^5}{\left(\frac{\pi}{4}\right) \cdot \rho \cdot T^3} = \frac{4}{\pi} \cdot \frac{Po \cdot N^3 \cdot D^5}{T^3}$$
(13)

Re-arranging for impeller speed:

$$N = \left(\frac{\pi \cdot \overline{\varepsilon} \cdot T^3}{4Po \cdot D^5}\right)^{1/3}$$
(14)

Substituting Equation (14) into Equation (12) gives the following:

$$\eta_{HYDR} = \frac{Fl}{Po \cdot D^2} \cdot \left(\frac{4}{\pi} \cdot \frac{4Po \cdot D^5}{\overline{\varepsilon} \cdot T^3}\right)^{2/3}$$

$$= 1.175 \frac{Fl}{Po^{2/3}} \cdot \left(\frac{D}{T}\right)^{4/3} \cdot \left(\overline{\varepsilon} \cdot T\right)^{-2/3}$$
(15)

The hydraulic efficiency data plotted in Figure 2 are replotted in Figure 3 using the new definition from Equation (15) with a power per mass of 1 W/kg and vessel diameter of 1 m. The effect of impeller diameter is now taken into account and large diameter impellers ($D/T \approx 0.5$) are more efficient than smaller ones ($D/T \approx 0.3$), pumping approximately twice the mass of fluid per unit of energy input.

Shear

 $\dot{\gamma} = \frac{v_H - v_L}{r_H - r_L}$

In any flowing system, the shear rate is the time-averaged velocity gradient [7].

Oldshue [3] has compared the time-averaged velocity gradients in the discharge of a hydrofoil and pitchedblade and Rushton turbines to show that the Rushton generates higher shear than the pitched-blade, which generates higher shear than the hydrofoil. This has become the conventional wisdom in the mixing field.

Figure 4 shows the mean velocity profiles for a hydrofoil (in green) and pitched-blade turbine (in red), which were measured using particle-image velocimetry in the PMSL laboratory. The dashed lines show the average velocity gradient in the discharge. Figure 5 shows the mean velocity profile for the Rushton turbine and, again, the dashed lines show the average velocity gradient in the discharge. The shear rate is described by the following equation:

FIGURE 3. Shown here is a plot of the hydraulic efficiency, η_{HYDR} , versus impeller-to-vessel-diameter ratio

where v_H and v_L are the high and low velocities in the gradient and r_H and r_L are the radial positions corresponding to the locations where these velocities were measured. Since the velocities are normalized by the impeller tip speed and the radial positions by the impeller radius, Equation (16) can be re-written as follows:

$$\dot{\gamma} = \frac{(\alpha - \beta) \cdot V_{TIP}}{(\psi - \omega) \cdot R} = \Lambda \cdot \frac{V_{TIP}}{R}$$
(17)

Values of α , β , ψ , ω and Λ are given in Table 2. Also the ratio of $\Lambda / \Lambda_{HYDFL}$ is shown and, at equal tip speed and impeller diameter the Rushton generates the highest shear rate followed by the pitched-blade turbine and then the hydrofoil.

Engineers are concerned with the power drawn by the impeller since this determines the size of the agitator needed to achieve the desired process result. Equation (13) can be rearranged to express the power input by the impeller per unit mass of fluid in terms of tip speed:

$$\overline{\varepsilon} = \frac{4}{\pi^4} \cdot Po \cdot \frac{V_{TIP}^3}{T} \left(\frac{D}{T}\right)^2$$
(18)

The π^3 term must be introduced because $V_{TIP} = \pi ND$.

Comparing different impellers of equal diameter at the same scale:

$$V_{TIP} \propto Po^{-1/3} \tag{19}$$

Comparing any impeller with the hydrofoil:

$$\frac{\dot{\gamma}_{IMP}}{\dot{\gamma}_{HYDFL}} = \frac{\Lambda_{IMP}}{\Lambda_{HYDFL}} \cdot \frac{V_{IMP}}{V_{HYDFL}}$$
(20)

Substituting Equation (19) into Equation (20):

$$\frac{\dot{\gamma}_{IMP}}{\dot{\gamma}_{HYDFL}} = \frac{\Lambda_{IMP}}{\Lambda_{HYDFL}} \cdot \left(\frac{Po_{HYDFL}}{Po_{IMP}}\right)^{1/3}$$
(21)

(16)

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TABLE 2: MEAN VELOCITY PROFILE SHEAR RATES									
Impeller	α	β	ψ	ω	Λ	$\Lambda/\Lambda_{\rm HYDFL}$ at equal $\rm V_{TIP}$	$\Lambda/\Lambda_{\rm HYDFL}$ at equal $\bar{\varepsilon}$		
Narrow-blade hydrofoil	0.25	0.14	0.80	0.16	0.17	1.00	1.00		
Pitched-blade turbine	0.35	0.07	0.65	0.12	0.53	3.11	1.91		
Rushton turbine	0.60	0.16	0.12	0.04	5.50	24.4	9.55		

Table 2 also shows the ratio of the shear rates when the impellers operate at equal power input per mass. Taking power numbers from Table 1, the ranking of the impellers does not change. Therefore, so, whether compared at equal tip speed or power input the Rushton turbine generates the highest shear rate followed by the pitched-blade turbine then the hydrofoil. This ranking can be tested against a process result that is dependent on shear, namely the break-up of droplets to create a liquidliquid dispersion.

Process result

Mass transfer between two immiscible liquid phases, with or without reaction, is an important process result. The interfacial area available for mass transfer is proportional to the volume fraction of the dispersed phase and inversely proportional to the Sauter mean droplet size [8].

If the dispersion is agitated for a long period of time (several hours), an "equilibrium droplet size" is achieved that is stable in the mixing environment in which the droplets are being formed. This means that there is an equilibrium between the forces breaking-up the droplets and the forces resisting break-up resulting from the interfacial tension between the two liquids and the viscosity of the dispersed phase liquid.

Figure 6 shows the Sauter mean droplet size plotted versus the average power input per unit mass for Rushton, two pitched-blade turbines, with blades angled at 45 and 60 deg, a hydrofoil and a high-shear disperser impeller. The experiments were carried out with low viscosity silicone oil as the dispersed phase and at a very low concentration so that the effect of coalescence on the droplet size can be ignored.

If the hydraulic efficiency and shear rate comparison quantify the performance characteristics of the impellers, when compared at the same power input per mass, the



FIGURE 4. This graph shows a plot of the mean velocity profiles for pitchedblade turbine and hydrofoil impellers

Rushton should create the smallest droplets and the hydrofoil, the largest, with the pitched-blade turbines falling somewhere between these two. In fact, the hydrofoil creates smaller droplets than the Rushton and two pitched-blade turbines and the droplets created by the turbines are indistinguishable experimentally. This result has also been observed by Pacek and others [9].

There is another geometrical property of impellers that determines how they create the "fluid dynamic effect" that achieves this desired "process result." This is the trailing vortices that form at the tip of the impeller blades.

Trailing vortex

As the impeller moves through the fluid, the pressure on the leading face of the blade is higher than on the back. The high- and low-pressure zones meet at the tip of the blade and the fluid moves from the high- to low-pressure region creating the trailing vortex. This phenomenon can often be observed on airplane wings [10, 11].

In a stirred tank, the velocity and size of the vortices can be measured using laser-Doppler or particle-image velocimetry, then the kinetic energy and energy dissipation rate (the local power input per mass) can be calculated. The kinetic energy of the trailing vortex is often non-dimensionalized by dividing by the impeller tip speed squared. Grenville and others [12] have shown that for impellers with blades:

$$\frac{\kappa_{MAX}}{V_{TIP}^2} = 0.104 \cdot Po^{1/2}$$
(22)

Where $V_{TIP} = \pi N \cdot D$.

The standard deviation for this correlation is $\pm 10\%$.

The maximum energy dissipation rate within the vortex is given by the following equation [13]:

$$\varepsilon_{MAX} = A \cdot \frac{k_{MAX}^{3/2}}{l_0} \tag{23}$$

 I_0 is a length scale related to the flow near the impeller and it is a fraction of the impeller diameter. Substituting Equation (22) into Equation (23) and setting $I_0 = D/x$ and A = 1:

$$\varepsilon_{MAX} = 1.04 \cdot x \cdot Po^{3/4} \cdot N^3 \cdot D^2 \tag{24}$$

The standard deviation for this correlation is $\pm 15\%$.

Where measurements have been made, typical values of *x* are given in Table 1. Again the value for the highshear disperser is similar to the Rushton, but there is no explanation as to why this is the case. Generally, the scale of the trailing vortex for the Rushton and pitchedblade turbines is equal to one-half of the projected height of the blade at its tip. For hydrofoils the scale of the trailing vortex is equal to the projected height of the blade at its tip.

Equations (13) and (24) can be combined to show that the ratio of the maximum energy dissipation rate to the average power input per mass, K, is:

$$K = \frac{\varepsilon_{MAX}}{\overline{\varepsilon}} = \frac{1.04 \cdot x \cdot Po^{3/4} \cdot N^3 \cdot D^2}{Po \cdot N^3 \cdot D^5} \cdot \frac{\pi T^3}{4}$$

= $0.82 \cdot \frac{x}{Po^{1/4}} \cdot \left(\frac{T}{D}\right)^3$ (25)

The ratio is weakly dependent on the type of impeller (*Po*), dependent on the scale of the vortex (x) and strongly dependent on the size of the impeller (*D*/*T*). The reason for this is that a small-diameter impeller must operate at a higher tip speed than a larger one to input the same power and the maximum kinetic energy is proportional to the tip speed squared.

The droplet size data plotted versus the average power input per mass in Figure 6 are replotted in Figure 7 versus the maximum energy dissipation rate in the trailing vortex. The variations in the trailing vortex energy dissipation rate generated by the impellers and the effects on the droplet size are now correctly accounted for, including the high-shear disperser.

The conventional wisdom in the mixing industry has been that hydrofoil impellers generate "low shear" and Rushton turbines generate "high shear" [3, 4] and this is true if only the time-averaged velocity gradients are compared. The maximum kinetic energy dissipation rate within the trailing vortex, ε_{MAX} , generates the stresses that break-up droplets, or any other second phase, in an agitated vessel. Rather than describing these impellers as "high shear," it is more rigorous to call them "high dissipation" or "high stress."

Applications

There are many processes in which the fluid dynamic effect that achieves the process result is commonly considered to be "shear" although, strictly, the process result is determined by the maximum energy dissipation rate within the trailing vortex. One example of a "shear" driven process is flocculation of fine particles. Agitators are designed to provide a desired shear rate, or *G*-value. *G* is defined as:

$$G = \left(\frac{P}{\mu \cdot V}\right)^{1/2} = \left(\frac{Po \cdot \rho \cdot N^3 \cdot D^5}{\mu \cdot V}\right)^{1/2}$$
(26)

This shear rate is based on the vessel-averaged power input per volume and the fluid's dynamic viscosity. Equation (26) suggests that, provided that the average power per volume is kept constant, the same *G*-value will be generated and the flocculation performance will be the same. Benz [14] has written a review of the problems that will be encountered taking this approach to agitator design, especially the fact that it takes no account of impeller type or diameter. He concludes that "*G*-value has no legitimate use in designing or specifying agitators."

Spicer and others [15] have measured the size and structure of flocculated polystyrene particles using a hydrofoil, pitched blade and Rushton turbines at G-values,



FIGURE 5. This graph shows a plot of the mean velocity profiles for the Rushton turbine

as defined in Equation (26), of 15, 25 and 50 s⁻¹. The corresponding values of vessel-averaged power input per mass are 2.25×10^{-4} , 6.25×10^{-4} and 2.50×10^{-3} W/kg. Grenville and Spicer [16] have re-analyzed these data and the floc length versus the maximum kinetic energy dissipation rate, calculated using Equation (23), is plotted in Figure 8. This approach to the analysis correlates the data and suggests that the concept of a *G*-value should work for agitator design provided that it is based on the maximum energy-dissipation rate in the trailing vortex — not the vessel averaged power per volume.

The selectivity of competitive reactions carried out in semi-batch mode is determined by the local mixing rate [17], the micro-mixing rate, in the region where the added reactant is introduced to the vessel [18]. Bourne and Dell'ava [19] have shown that the selectivity of an azo-coupling reaction can be maximized by feeding the added reactant at the impeller where the trailing-vortex energy dissipation rate determines the rate of micro-mixing. They, and Nienow and others [20], have also shown that, provided the feed location is geometrically similar, the selectivity of the reaction can be maintained on scaleup if the trailing-vortex energy dissipation rate is the same at the two scales. This has also been shown to apply to precipitation reactions where the particle size and morphology need to be controlled [21, 22].



FIGURE 6. This graph shows the behavior of the Sauter mean diameter, d_{32} , versus vessel-averaged power input per unit mass



FIGURE 7. Shown here is a plot of the Sauter mean diameter, d_{32} , versus the trailing vortex energydissipation rate

Finally, mixing in crystallization processes requires both rapid local mixing to minimize primary nucleation and high flow to promote homogeneity, favoring secondary nucleation and crystal growth. Also, a balance between crystal growth and crystal damage must be considered in choosing the appropriate impeller [23].

Conclusions

Mixing processes can be described in terms of the desired process result. Generally this result will be controlled by the flow and turbulence intensity generated by an impeller. The approach described here can be used to determine which the best impeller to achieve this result will be. It can also be used to translate laboratory and pilot-scale results taken with one type of impeller to a larger scale using a different geometry, provided that the process result and controlling dynamic effect can be identified.

The term high-shear is commonly used to describe an impeller's capability for dispersion of a second immiscible phase generating surface area for mass transfer. Similarly, low-shear is used to describe impellers that, in multi-phase processes, allow the second phase to grow, and flocculation is a good ex-



FIGURE 8. This plot shows the maximum average floc length versus maximum kinetic-energy-dissipation rate for hydrofoil, pitched blade and Rushton turbines

ample of this.

In a turbulent agitated vessel, the time-averaged velocity gradients are of little use, and potentially misleading, for comparison of impeller performance and agitator design. While the term "shear" is used qualitatively to describe impellers' dispersing capabilities, it must be recognized that the true mechanism of break-up is determined by the maximum energy dissipation rate within the impellers' trailing vortices. This understanding enables engineers to select the appropriate impellers for their processes.

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